

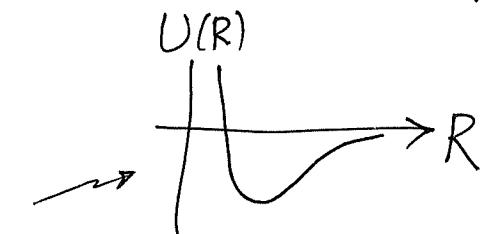
D. Born-Oppenheimer Approximation

[[↑]Max Born] [[↑]J. Robert Oppenheimer]

- Made approximations to justify separation (decoupling) of Ψ_{molecule} into electronic part and nuclear part

$$\Psi_{\text{molecule}}(\vec{R}; \vec{r}_1, \dots, \vec{r}_N) \approx \psi_N(\vec{R}) \cdot \psi_{\text{el}}(\vec{r}_1, \vec{r}_2, \dots, \vec{r}_N; \vec{R}) \quad \begin{matrix} \text{nuclei fixed} \\ \text{with separation } R \end{matrix} \quad (7)$$

- Why is it a good approximation?
- How bonding comes out from solving ψ_{el} ?
- How vibrational and rotational motions come out from $\psi_N(\vec{R})$?



Key idea

Key idea 
 $m \ll \mu$ $\frac{m}{\mu} \sim 10^{-3}$ to $10^{-5} \ll 1$

time scale of nucleus motion \gg time scale of electrons' motion
[slow] [fast]

⇒ for electrons, in their time scale, they "see" a potential energy function due to the nuclei as if the nuclei are fixed with a separation R

⇒ Solve Schrödinger Equation for electrons only with nuclei separated by R (with R being a parameter, not an operator any more) (8)

Meaning:

nuclei k.e. (for electronic problem, ignore for now)

$$\hat{H}_{\text{total}} \approx \left(-\frac{\hbar^2}{2\mu} \nabla_R^2 \right) + \left(\sum_{i=1}^N \left(-\frac{\hbar^2}{2m} \nabla_{\vec{r}_i}^2 \right) - \sum_{i=1}^N \left(\frac{Z_A e^2}{4\pi\epsilon_0 |\vec{r}_i - \vec{R}_A|} + \frac{Z_B e^2}{4\pi\epsilon_0 |\vec{r}_i - \vec{R}_B|} \right) + \frac{Z_A Z_B e^2}{4\pi\epsilon_0 R} \right) + \sum_{i < j} \frac{e^2}{4\pi\epsilon_0 |\vec{r}_i - \vec{r}_j|}$$

involves electrons $\{\vec{p}_i, \vec{r}_i\}$ only

$$= \left(-\frac{\hbar^2}{2\mu} \nabla_R^2 \right) + [\hat{H}_{\text{electronic}} (\vec{p}_1, \dots, \vec{p}_N; \vec{r}_1, \dots, \vec{r}_N; R)] \quad (9)$$

Consider this
only after solving

\hat{H}_{el} Schrödinger problem

involves only electronic variables (with R being a constant)

[Important: Will be back to Eq.(8) later]

Aside: For those who like math...

[downgraded from operator]
treated as a constant
for now

$$\hat{H} \approx -\frac{\hbar^2}{2M} \nabla_{\vec{R}}^2 + \hat{H}_{\text{electronic}}(\vec{p}_1, \dots, \vec{p}_N; \vec{r}_1, \dots, \vec{r}_N; R)$$

nuclear separation
 Nuclear Variables Only electrons' variables Only

Separation of Variables using Eq.(9) and

$$\Psi_{\text{molecule}}(\vec{R}; \vec{r}_1, \dots, \vec{r}_N) = \psi_N(\vec{R}) \cdot \psi_{el}(\vec{r}_1, \vec{r}_2, \dots, \vec{r}_N; R)$$

⇒ an equation for $\psi_{el}(\vec{r}_1, \vec{r}_2, \dots, \vec{r}_N; R)$ (bonding) [see Eq.(10)]

an equation for $\psi_N(\vec{R})$ (vibrational, rotational effects)

2-step process to solve Eq.(9)

Step 1: Solve electronic part (with R as parameter)

$$\hat{H}_{\text{electronic}} \psi_{\text{el}}(\vec{r}_1, \dots, \vec{r}_N; R) = E_{\text{electronic}}(R) \psi_{\text{el}}(\vec{r}_1, \dots, \vec{r}_N; R) \quad (10)$$

- Many-electron QM problem (c.f. multi-electron atoms)

- Which R to use? [Try different R 's. One R is one problem to solve!]

Key idea

$$\bullet \leftarrow R_1 \rightarrow \bullet$$

Eq.(10) gives $E_{\text{el}}(R_1)$

$$\bullet \leftarrow R_2 \rightarrow \bullet$$

Eq.(10) gives $E_{\text{el}}(R_2)$

$$\bullet \leftarrow R_3 \rightarrow \bullet$$

Eq.(10) gives $E_{\text{el}}(R_3)$

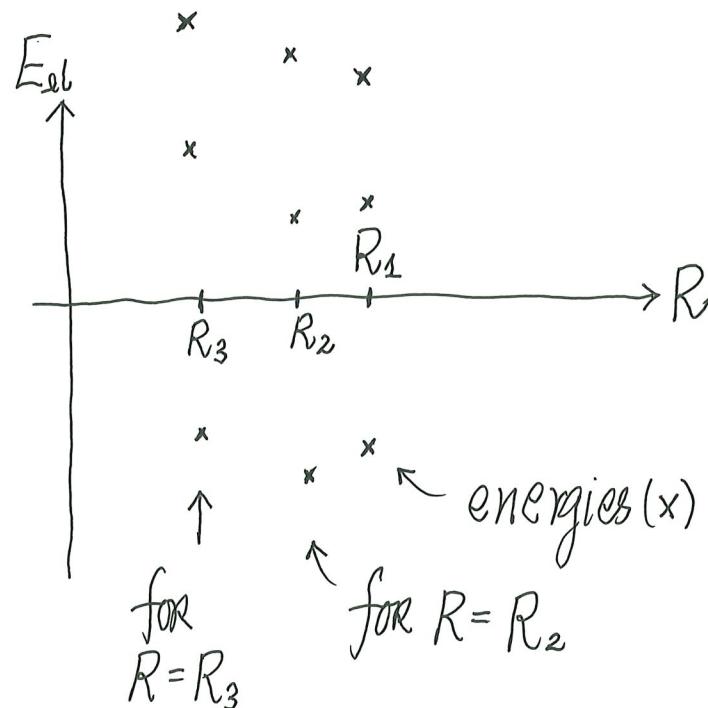
and

so

on

⋮
⋮
⋮

Each is a different
QM problem for
solving Eq.(10)



$$\hat{H}_{\text{electronic}} \psi_{\text{el}}(\vec{r}_1, \dots, \vec{r}_N; R) = E_{\text{el}} \psi_{\text{el}}(\vec{r}_1, \dots, \vec{r}_N; R) \quad (10)$$

↑
labels value of R used in solving TISE⁺

energies (x) solving Eq.(10) for $R = R_1$

for $R = R_3$
↑
for $R = R_2$

Recall: Solving TISE gives
many energy eigenvalues

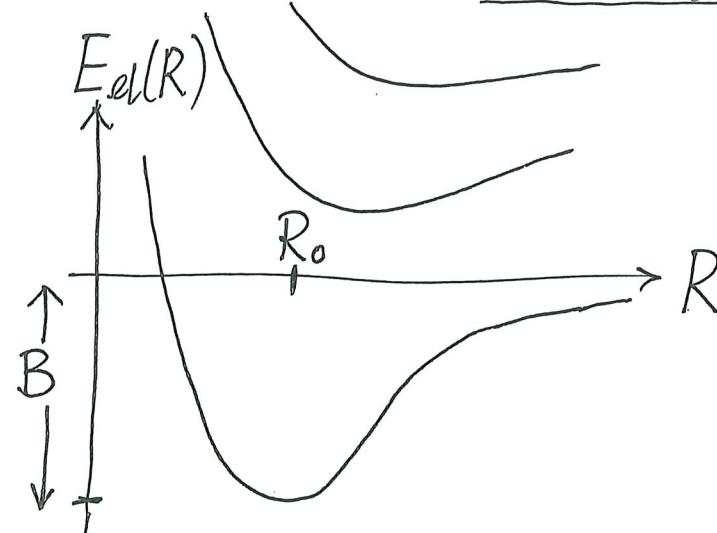
Different $\hat{H}_{\text{electronic}}$ \Rightarrow Different TISE's

Repeat calculations for many different inter-nucleus separations R
 \Rightarrow many curves $E_{\text{el}}(R)$

⁺ Eq.(10) is a many-electron TISE just like atomic problems. How can we "solve" it?
As in atomic problems, we need to make further approximations.

- At the end, we have how E_{el} depends on R , i.e. $E_{el}(R)$

Possibly,



shows preferred separation
bond length R_0
and Binding energy B

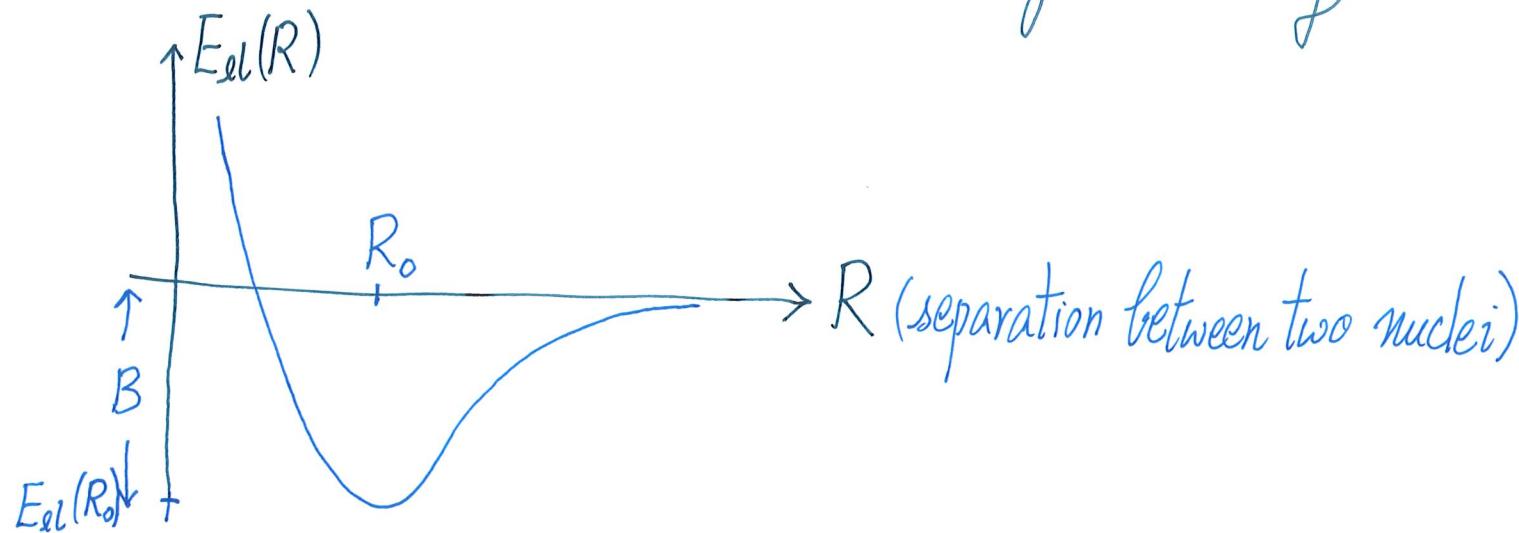
∴ Electronic problem gives the Quantum Physics of
Chemical Bonds!

Q: Why does the CO (carbon monoxide) molecule have bond length 1.128\AA ?

We have $\hat{H}_{\text{electronic}}^{(\text{CO})} \rightarrow \text{TISE} \rightarrow E_{el}(R) \rightarrow R_0$

electrons behavior ($|\psi_{el}|^2$) that selects preferred separation R_0

Due to the electrons, there emerges binding,



If nuclear motion is ignored, electrons [their distribution] select a certain bond length R_0 and certain bond strength B



and energy of the molecule would be $E_{el}(R_0)$

The picture of the molecule up to here is [not true]



[no vibration, no rotation YET!]

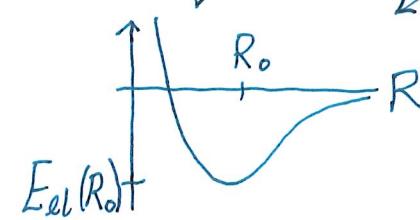
Step 2 : Go back to \hat{H}_{total} , and consider nuclear motion

with $-\frac{\hbar^2}{2\mu} \nabla_R^2$ term

new physics in molecules

At this point, the electrons are done with their duties.

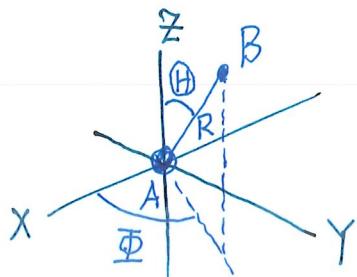
- What's left is the relative motion problem.



3D One-body of mass μ under the influence of $E_{\text{el}}(R)$

depends on separation R only (spherically symmetric)
[Generally R, θ, ϕ of mass μ]

Picture



- rod could rotate
- could vibrate (radial direction)

capital symbols used for nucleus B's coordinates

The situation has a 2D analogy (easier to visualize)

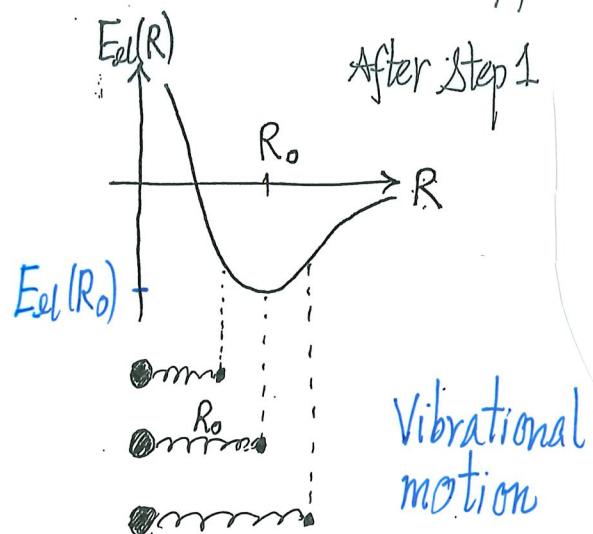
Radial motion is approximately harmonic oscillator

Near minimum,

$$E_{el}(R \approx R_0) = E_{el}(R_0) + \frac{1}{2} \left(\frac{d^2 E_{el}}{dR^2} \right)_{E_0} (R - R_0)^2$$

$$= E_{el}(R_0) + \frac{1}{2} k (R - R_0)^2$$

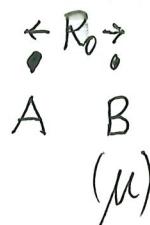
$$k \equiv \frac{1}{2} \left(\frac{d^2 E_{el}}{dR^2} \right)_{E_0} \quad (\text{due to electrons})$$



R_0 = equilibrium separation (bond length)
= "natural length" of spring (bond)

"Spring Constant"?

"Curvature near R_0 "

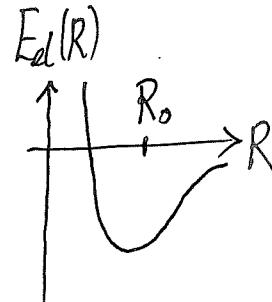


stiffness of the chemical bond
QM of harmonic oscillator

allowed energies are

$$(n + \frac{1}{2})\hbar\omega$$

$$= (n + \frac{1}{2}) \hbar \sqrt{\frac{k}{\mu}} \quad (11)$$

2D analogy

rotational motion
at the minimum
(assume $R=R_0$)

Angular Motion

rotational motion

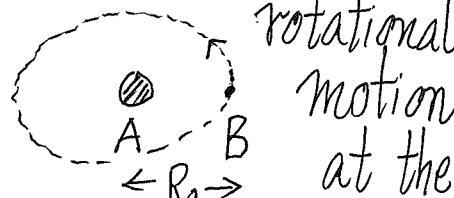
"2D rigid rotor"

means $R=R_0$

Allowed energies are

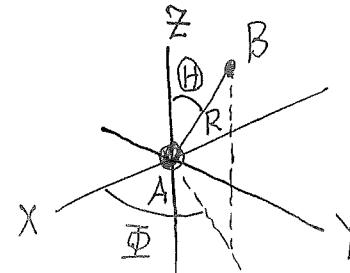
$$\frac{L^2}{2I} = \frac{L^2}{2\mu R_0^2}$$

moment of inertia

Diatomic Molecules

Angular motion is rotational motion

"3D rigid rotor"



B freely moves on
sphere of radius R_0

Allowed energies are

$$\frac{l(l+1)\hbar^2}{2\mu R_0^2} \quad (12)^+$$

"rotational levels"

[Chemists use $\frac{J(J+1)\hbar^2}{2\mu R_0^2}$]

At the end (Steps 1+2), an allowed energy of a molecule is specified by

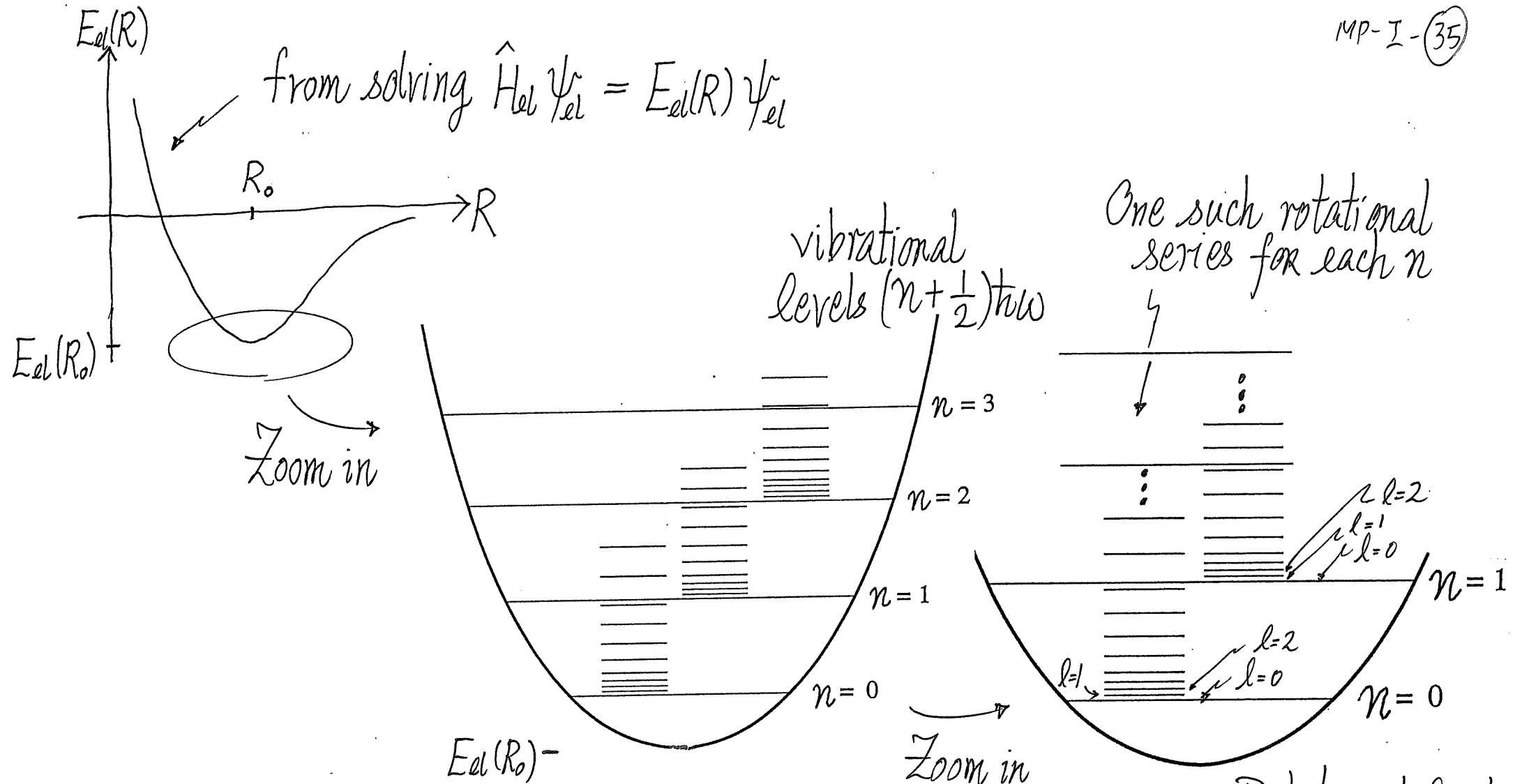
$$\boxed{E_{\text{total}}^{(\text{molecule})} = \underbrace{E_{\text{el}}(R_0)}_{\text{electronic part}} + \underbrace{(n + \frac{1}{2})\hbar\omega}_{\text{vibrational part}} + \underbrace{\frac{l(l+1)\hbar^2}{2\mu R_0^2}}_{\text{rotational part}}}$$

(13)

[Key result without going through the mathematics!]

①
↑
electrons bind the atoms at set the scale of the binding $E_{\text{el}}(R_0)$, the spring constant k (thus ω) and equilibrium bond length R_0 for...

②
↑
... the vibrational motion of the nuclei with $\omega = \sqrt{\frac{k}{\mu}}$...
... and the rotational motion of the nuclei



An energy diagram showing the rotational levels associated with each vibrational state in diatomic molecule.

Picture of Key Result Eq.(13)

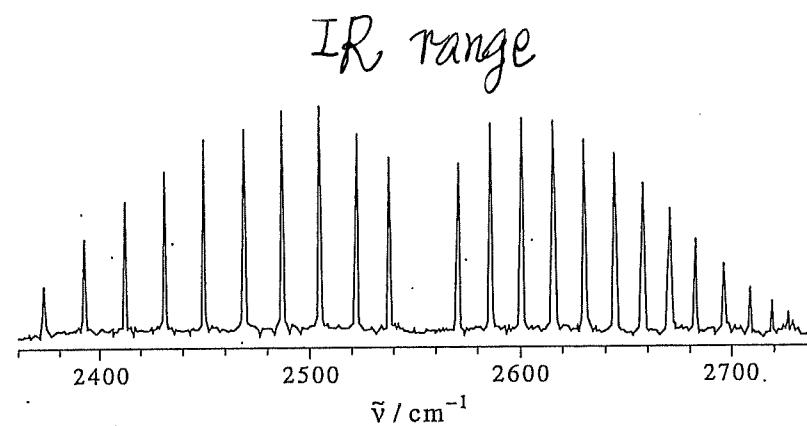
Rotational levels

$$\frac{l(l+1)\hbar^2}{2I}$$

 [each l , $(2l+1)$ values of m]
 degeneracy

Experimental consequences

- Transitions between one molecular state $[n, l]$ to another state $[n', l']$ give finger-print molecular spectrum of the molecule



Rotational-vibrational Spectrum of HBr
 [Questions: Get R_0 and $\omega = \sqrt{k/m}$ out from spectrum?]

- Must understand the physical picture behind Born-Oppenheimer approximation first, before going into more detail
- It is important also because the same physics works in understanding solids. ["Solid is just a huge molecule"]
- Eq.(13) is the key result for understanding molecular spectrum [c.f. atomic spectra motivated the development of QM]
- This ends the discussion on the Big Picture of the QM theory of molecular states